

Schwartz 6.2

Taking x_1 as source, the retarded ~~into~~ propagator is

$$D(x_1, x_2) = \langle 0 | \phi_0(x_2) \phi_0(x_1) | 0 \rangle$$

$$= \langle 0 | \int \frac{d^3 \vec{k}_1}{(2\pi)^3} \frac{d^3 \vec{k}_2}{(2\pi)^3} [a_{\vec{k}_2} e^{-i\vec{k}_2 x_2} + a_{\vec{k}_2}^\dagger e^{i\vec{k}_2 x_2}]$$

$$\times [a_{\vec{k}_1} e^{-i\vec{k}_1 x_1} + a_{\vec{k}_1}^\dagger e^{i\vec{k}_1 x_1}] | 0 \rangle \frac{1}{\sqrt{\omega_{\vec{k}_1}}} \frac{1}{\sqrt{\omega_{\vec{k}_2}}}$$

$$= \int \frac{d^3 \vec{k}_1}{(2\pi)^3} \frac{d^3 \vec{k}_2}{(2\pi)^3} \langle 0 | a_{\vec{k}_2} a_{\vec{k}_1}^\dagger | 0 \rangle e^{-i\vec{k}_2 x_2} e^{i\vec{k}_1 x_1} \frac{1}{\sqrt{\omega_{\vec{k}_1}}} \frac{1}{\sqrt{\omega_{\vec{k}_2}}}$$

$$= \int \frac{d^3 \vec{k}_1}{(2\pi)^3} \frac{d^3 \vec{k}_2}{(2\pi)^3} (2\pi)^3 \delta^3(\vec{k}_1 - \vec{k}_2) \frac{1}{\sqrt{\omega_{\vec{k}_1}}} \frac{1}{\sqrt{\omega_{\vec{k}_2}}} e^{-i\vec{k}_2 x_2} e^{i\vec{k}_1 x_1}$$

$$= \int \frac{d^3 \vec{k}}{(2\pi)^3} \frac{1}{2\omega_{\vec{k}}} e^{-i\vec{k}(x_2 - x_1)}$$

$$= \int \frac{d^4 k}{(2\pi)^3} \frac{1}{2k_0} e^{-i\vec{k}(x_2 - x_1)} f(k_0 - |\vec{k}|)$$

↑
(propagating field have $m=0$)

Similarly, the advanced propagator would be

$$\int \frac{d^4 k}{(2\pi)^3} \frac{1}{2k_0} e^{-i\vec{k}(x_1 - x_2)} f(k_0 - |\vec{k}|)$$

The difference is that ~~the~~ ~~retard~~ we have $(x_2)_0 \geq (x_1)_0$ for retarded, and $(x_1)_0 \geq (x_2)_0$ for advanced propagator.